



# Improved Numerical Generalization of the Bethe-Weizsäcker Mass Formula for Prediction the Isotope Nuclear Mass, the Mass Excess Including of Artificial Elements 119 and 120

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<sup>\*</sup>Corresponding author**To cite this article:**Mavrodiev Strachimir Chterev, Vol Alexander. Improved Numerical Generalization of the Bethe-Weizsäcker Mass Formula for Prediction the Isotope Nuclear Mass, the Mass Excess Including of Artificial Elements 119 and 120. *Nuclear Science*. Vol. 4, No. 2, 2019, pp. 11-22.

doi: 10.11648/j.ns.20190402.11

**Received:** June 5, 2019; **Accepted:** July 15, 2019; **Published:** July 26, 2019

**Abstract:** George Gamow's liquid drop model of the nucleus can account for most of the terms in the formula and gives rough estimates for the values of the coefficients. Its semi-numerical equation was first formulated in 1935 by Weizsäcker and in 1936 Bethe [1, 2], and although refinements have been made to the coefficients over the years, the structure of the formula remains the same today. Their formula gives a good approximation for atomic masses and several other effects, but does not explain the appearance of magic numbers of protons and neutrons, and the extra binding-energy and measure of stability that are associated with these numbers of nucleons. Mavrodiev and Deliyergiyev [3] formalized the nuclear mass problem in the inverse problem framework. This approach allowed them to infer the underlying model parameters from experimental observation, rather than to predict the observations from the model parameters. They formulated the inverse problem for the numerically generalized semi-empirical mass formula of Bethe and von Weizsäcker going step-by-step through the AME2012 [4] nuclear database. The resulting parameterization described the measured nuclear masses of 2564 isotopes with a maximal deviation of less than 2.6 MeV, starting from the number of protons and number of neutrons equal to 1. The unknown functions in the generalized mass formula was discovered in a step-by-step way using the modified procedure realized in the algorithms developed by Aleksandrov [5-7] to solve nonlinear systems of equations via the Gauss-Newton method. In the presented herein article we describe a further development of the obtained by [3] formula by including additional factors,- magic numbers of protons, neutrons and electrons. This inclusion is based the well-known experimental data on the chemically induced polarization of nuclei and the effect of such this polarization on the rate of isotope decay. It allowed taking into account resonant interaction of the spins of nuclei and electron shells. As a result the maximal deviation from the measured nuclear masses of less than 1.9 MeV was reached. This improvement allowed prediction of the nuclear characteristics of the artificial elements 119 and 120.

**Keywords:** Bethe-Weizsäcker Mass Formula, Magic Numbers, Binding Energy, Wigner Term, Inverse Problem, Electrons-Nucleus Interaction, Chemical Polarization, Isotopes

## 1. Introduction

In the last few years there has appeared new experimental data which demonstrated the dramatic change of decay rate due to the ionization of an atom and due to the resonant interaction between the electron shells and the nuclei [8-14].

For example, a strong dependence of the nuclear decay rate on ionization was shown for the  $^{229}\text{Th}$  90,  $^{226}\text{Rn}$  88,  $^{152}\text{Eu}$  63,  $^{154}\text{Eu}$  63 isotopes and the  $^{178}\text{mHf}$  72,  $^{99}\text{mTc}$  43 isomers. Testing the effectiveness of accounting for the

interaction of nuclei and the electron shell for accuracy of a well-known formula expands our understanding of the structure of atoms and the possible contribution of nucleus interaction to the formation of compounds and biological structures. Deepening our understanding of the nature of the interaction of nuclear and chemical processes can be very important both in solving the problem of radioactive waste and in solving a number of problems in biology and medicine [15].

Following are the steps involved in the generalization of the BW mass formula.

## 2. Bethe-Weizsäcker Mass Formula Digital Generalization

### 2.1. Original Bethe-Weizsäcker Mass Formula [1, 2]

Let A be the total number of nucleons, Z the number of protons, and N the number of neutrons, so that  $A = Z + N$ . The mass of an atomic nucleus will be

$$m = Z \cdot m_p + N \cdot m_n - B_E,$$

### 2.2. Hypothesis for Digital Generalization of Bethe-Weizsäcker Mass Formula with the Influence of Magic Numbers [3, 16-20]

$$B_E(Z, N, a) = Vol(Z, N, a) - Sur(Z, N, a) \frac{1}{A^{P_1(Z, N, a)}} - Cha(A, Z, a) \frac{Z(Z-1)}{A^{P_2(Z, N, a)}} - Sym(Z, N, a) \frac{(N-Z)^2}{A^{P_3(Z < N, a)}} + Wig(Z, N, a) \frac{\delta(N, Z)}{A^{P_4(Z, N, a)}} + K_{MN}(Z, N, a),$$

Where the function  $K_{MN}(Z, N, a)$  depends on the proton and neutron magic numbers and where the frontier between their influence and a is a set of unknown digital parameters. The function  $\delta$  is defined as:

$$\delta(N, Z) = +1 \text{ for even } N, Z$$

$$\delta(Z, N) = -1 \text{ for odd } N, Z \text{ and}$$

$$\delta(Z, N) = 0 \text{ for odd } A = Z + N$$

$$B_E(Z, N, a) = Vol(Z, N, a) - Sur(Z, N, a) \frac{1}{A^{P_1(Z, N, a)}} - Cha(A, Z, a) \frac{Z(Z-1)}{A^{P_2(Z, N, a)}} - Sym(Z, N, a) \frac{(N-Z)^2}{A^{P_3(Z < N, a)}} + Wig(Z, N, a) \frac{\delta(N, Z)}{A^{P_4(Z, N, a)}} + \frac{SVol(N, Z, E)}{A^{P_5(Z, N, a)}} + K_{MN}(Z, N, a) \quad (4)$$

Where: E is the number of electrons in a shell.

Eight electron magic numbers (2, 10, 18, 36, 54, 86, 118, 140) - from the periodic Mendeleev table of elements.

The method of discovering the explicit form of the functions  $Vol(Z, N, a)$ ,  $Sur(Z, N, a)$ ,  $Cha(A, Z, a)$ ,  $Sym(Z, N, a)$ ,  $Wig(Z, N, a)$ ,  $SVol(N, Z, E)$ ,  $K_{MN}(Z, N, a)$ ,  $P_1(Z, N, a)$ ,  $P_2(Z, N, a)$ ,  $P_3(Z, N, a)$ ,  $P_4(Z, N, a)$ ,  $P_5(Z, N, a)$  and the calculated values of a set of unknown digital parameters is described in papers [3, 5-7, 16-20] using the 3 overdetermined algebraic system of the equality of experimental values of equations (1, 2, 3) and their model,

$$V1 = \frac{Z}{A}, V2 = V12, V3 = V13, V4 = \frac{N}{A}, V5 = V42, V6 = V43, V7 = \frac{N-Z}{A}, V8 = V72, V9 = V73,$$

$\{ \text{displaystyle } m = Z \cdot m_p + N \cdot m_n - \frac{E_B}{A} \} \}$  where  $m_p$  and  $m_n$  are the rest mass of proton and neutron, respectively, and  $E_B$  is the binding energy of the nucleus. The semi-empirical mass formula states that the binding energy  $E_B$  per nucleon will take the following form:

$$BE = Vol - Sur - Cha - Sym + Wig.$$

In this paper we use the connections between atomic masses, nuclear masses, and mass excess as follows:

$$AtMass = Z \cdot mp + N \cdot mn - BE \quad (1)$$

$$NuclMass = AtMass - (mE + Ael E2.39 + Bel E5.35) \quad (2)$$

$$MassExc = AtMass - Au \quad (3)$$

where  $mn = 939.565301$ ,  $mp = 938.2719982$ ,  $ME = 0.510998902$ ,  $1u = 931.494028 \text{ MeV}/c^2$  (2006, CODATA) correspondingly with fit parameters

$Ael = 1.44381E-05$ ,  $Bel = 1.554680E-12$ , where E is the number of electrons in a shell.

### 2.3 Hypothesis for Digital Generalization of Bethe-Weizsäcker Mass Formula with the Influence of the Magic Numbers and the Influence of the Electron Shell

Here  $B_E(Z, N, a)$  Is Defined as

calculated from equation (4). The number of the experimental data from the data base [21, 22] is M = 2536.

### 2.4. About the Choice of Arguments for Solving the Inverse Problem

Concerning the possibilities of REGN code, it is very convenient to choose the arguments which are linearly independent as well as with a variation near to zero. In our case we chose the arguments as follow:

$$\begin{aligned}
 V10 &= \frac{Z}{N+1}, V11 = V102, V12 = V103, V13 = \frac{1}{\ln(A+1)}, V14 = \frac{Z-Z_0}{Z+Z_0}, V15 = \frac{N-N_0}{N+N_0}, V16 = \frac{E-E_0}{E+E_0}, \\
 V17 &= V142, V18 = V152, V19 = V162, V20 = \frac{1}{A} \text{ if } A \text{ is even, and } -\frac{1}{A} \text{ if } A \text{ is odd,} \\
 V21 &= \frac{\pm 1}{Z+1} \text{ and } V22 = \frac{\pm 1}{N+1} \text{ correspondingly.}
 \end{aligned}$$

### 3. The Explicit Form of Unknown Functions

These forms are as follows:

$$\begin{aligned}
 Vol(Z, N, a, i) &= \exp(a_1 + \text{CorPow}(a, N, \text{isp})) + \text{CorS}(a, N, N_0 + 6), \\
 Sur(Z, N, a, i) &= \exp(a_2 + \text{CorPow}(a, N, \text{isp} + n\text{Pow})) + \text{CorS}(a, N, N_0 + 12), \\
 Cha(Z, N, a, i) &= \exp(a_3 + \text{CorPow}(a, N, \text{isp} + 2 n\text{Pow})) + \text{CorS}(a, N, N_0 + 18), \\
 Sym(Z, N, a, i) &= \exp(a_4 + \text{CorPow}(a, N, \text{isp} + 3 n\text{Pow})) + \text{CorS}(a, N, N_0 + 24), \\
 Wig(Z, N, a, i) &= \exp(a_5 + \text{CorPow}(a, N, \text{isp} + 4 n\text{Pow})) + \text{CorS}(a, N, N_0 + 30), \\
 SVol(Z, N, a, i) &= a_{N_0 + 6}(\exp(a_6 + \text{CorPow}(a, N, \text{isp} + 5 n\text{Pow}))) + \text{CorS}(a, N, N_0 + 36) v_{19}, \\
 P_1(Z, N, a, i) &= \exp(a_7 + \text{CorPow}(a, N, \text{isp} + 6 n\text{Pow})) + \text{CorS}(a, N, N_0 + 42), \\
 P_2(Z, N, a, i) &= \exp(a_8 + \text{CorPow}(a, N, \text{isp} + 7 n\text{Pow})) + \text{CorS}(a, N, N_0 + 48), \\
 P_3(Z, N, a, i) &= \exp(a_9 + \text{CorPow}(a, N, \text{isp} + 8 n\text{Pow})) + \text{CorS}(a, N, N_0 + 54), \\
 P_4(Z, N, a, i) &= \exp(a_{10} + \text{CorPow}(a, N, \text{isp} + 9 n\text{Pow})) + \text{CorS}(a, N, N_0 + 60), \\
 P_5(Z, N, a, i) &= \exp(a_{11} + \text{CorPow}(a, N, \text{isp} + 10 n\text{Pow})) + \text{CorS}(a, N, N_0 + 66),
 \end{aligned}$$

Where:

$$\begin{aligned}
 \text{CorPow}(a, N, i) &= \exp(-(\sum_{j=1}^{19} a_{i+j} v_j)^2), \\
 \text{CorS}(a, N, i) &= \exp(-(a_{i+1} v_{14} + a_{i+2} v_{15} + a_{i+3} v_{16} + a_{i+4} v_{20} + a_{i+5} v_{21} + a_{i+6} v_{22})^2)
 \end{aligned}$$

and

$$K_{MN}(Z, Z_0, WZ, N, N_0, WN, a, i) = K(Z, Z_0, WZ, N, N_0, WN, a, \text{isp} + 11 n\text{Pow})(1 + a_{N_0 + 5} v_3),$$

Where:

$$\begin{aligned}
 K(Z, Z_0, WZ, N, N_0, WN, a, \text{isp} + 11 n\text{Pow}) &= BZ(Z, Z_0, WZ, a, \text{isp} + 11 n\text{Pow})(1 + \exp(-(a_{N_0 + 1} v_1)^2 + a_{N_0 + 3} v_1) \\
 &+ BN(N, N_0, WN, a, \text{isp} + 11 n\text{Pow})(1 + \exp(-(a_{N_0 + 2} v_2)^2 + a_{N_0 + 4} v_4)), \\
 BZ(Z, Z_0, WZ, N, N_0, WN, a, i) &= Az(Z, Z_0, N, N_0, a, i) \exp(-\frac{(Z-Z_0)^2}{Gz(Z, Z_0, WZ, N, N_0, a, i)}) / ((Z - Z_0)^2 + Gz(Z, Z_0, WZ, N, N_0, a, i)), \\
 BN(Z, Z_0, WZ, N, N_0, WN, a, i) &= An(Z, Z_0, WZ, N, N_0, a, i) \exp(-\frac{(N-N_0)^2}{Gn(Z, Z_0, N, N_0, WN, a, i)}) / \\
 &((N - N_0)^2 + Gn(Z, Z_0, N, N_0, WN, a, i)), \\
 Az(Z, Z_0, N, N_0, a, i) &= \text{CorAmp}(a, n, i), An(Z, Z_0, N, N_0, a, i) = \text{CorAmp}(a, n, i + n\text{BWA}), \\
 Gz(Z, Z_0, WZ, N, N_0, a, i) &= WZ + \text{CorAmp}(a, n, i + 2 n\text{BWA}), \\
 Gn(Z, Z_0, N, N_0, WN, a, i) &= WN + \text{CorAmp}(a, n, i + 3 n\text{BWA}), \\
 \text{CorAmp}(a, n, i) &= \exp(a_{i+20} - (\sum_{j=1}^{19} a_{i+j} v_j)^2),
 \end{aligned}$$

Where the values of integer numbers are as follows:

iStr = 6, iPow = 5, iSP = iStr + iPow, nPow = 19, nBWA = 20, nBW = 4 nBWA, MnZ = 10, MnN = 11, Dop = 72, N0 = iSP(1 +

$$n\text{Pow}) + n\text{BW}, N = N_0 + \text{Dop} + 2 \text{MnZ} + 2 \text{MnN}.$$

$Z_0$  is the nearest of a set of proton magic numbers: 2, 8, 14, 20, 28, 50, 82, 98, 108, 124,

$N_0$  is the nearest of a set of neutron magic numbers: 2, 8, 14, 20, 28, 50, 82, 124, 152, 184, 202.

The values of magic numbers  $Z_0$ ,  $N_0$  as well as the boundaries between them  $WN$ ,  $WN$ , (half sum of consequently magic numbers) are the result of a fit procedure.

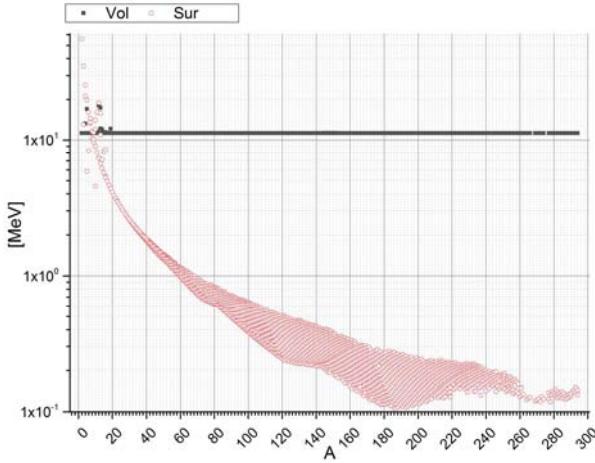


Figure 1. The behaviour of volume and surface terms.

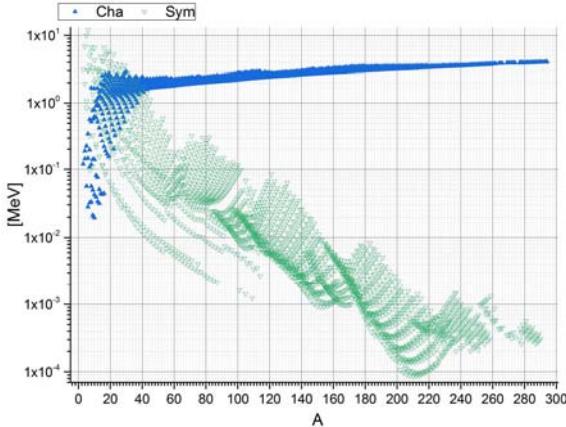


Figure 2. The behaviour of charge and symmetry terms.

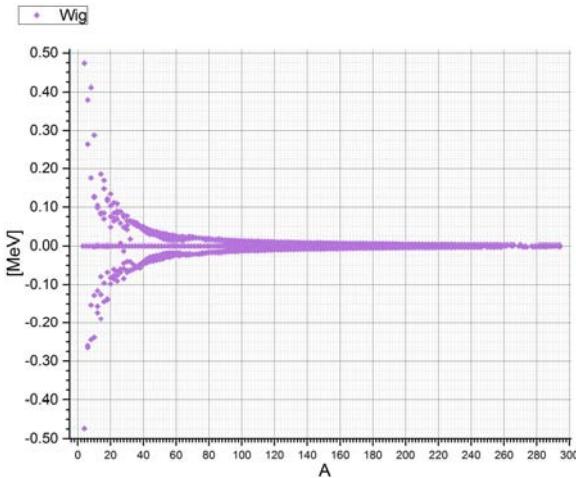


Figure 3. The behaviour of Wigner term.

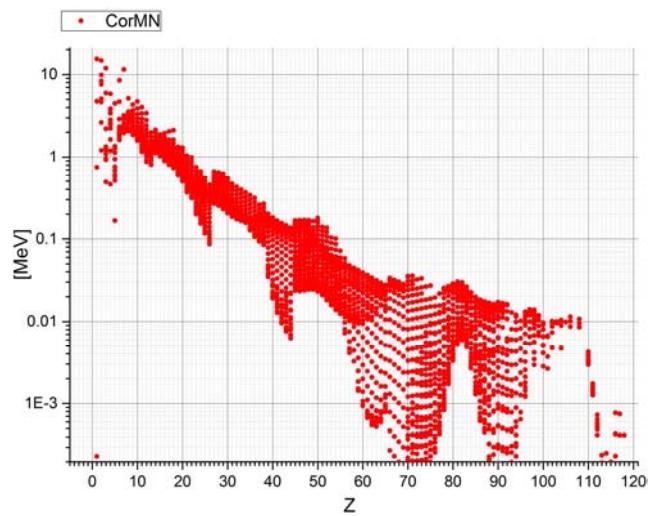


Figure 4. The behaviour of proton magic number influence.

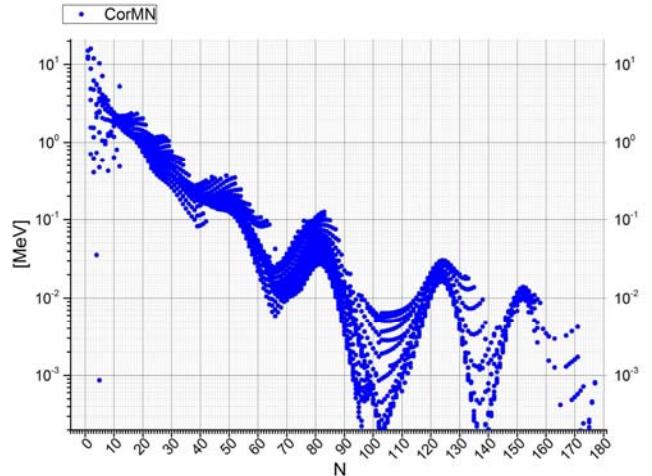


Figure 5. The behaviour of neutron magic number influence.

## 4. Description of Data

The following figures are a description of the data calculated.

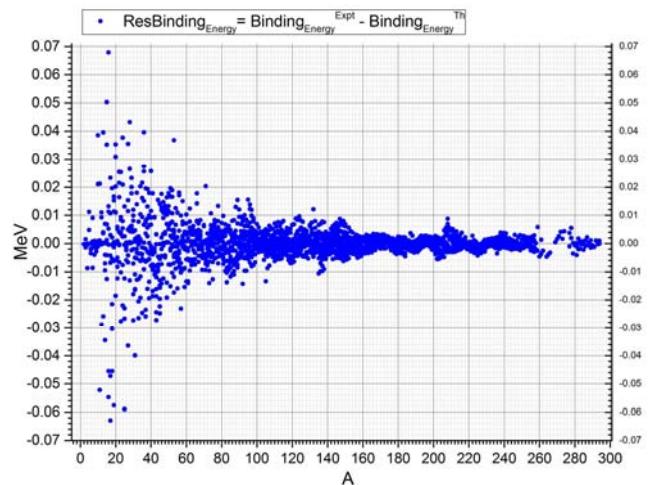


Figure 6. The behaviour of binding energy residuals.

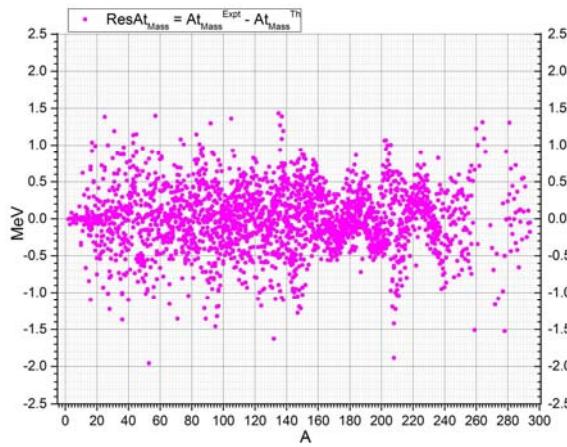


Figure 7. The behaviour of atomic masses residuals.

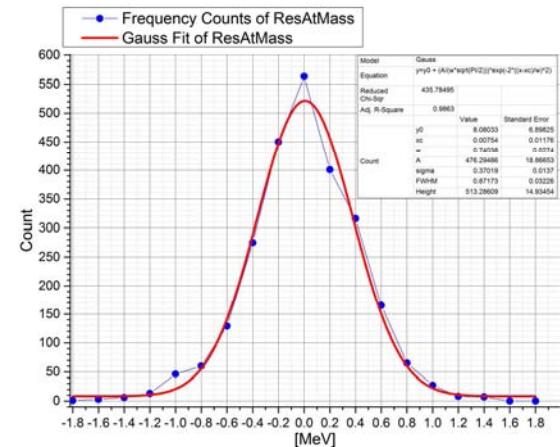
Figure 8. Gauss fit of atomic masses residuals ( $w = 0.740 \pm 0.027 \text{ MeV}$ ).

Table 1. Illustrates the quality of description for Binding Energy, Nuclear Mass, Mass Excess and Atomic Masses as function of Z, N and A.

Parameter	N total	Mean	SD	Variance	RMS	MAD	Minimum	Median	Maximum
Residual Binding energy	2536	-2.54E-05	0.00688	4.74E-05	0.008	0.00369	-0.06291	3.49E-06	0.06791
Residual Nuclear mass	2536	-1.03E-05	0.41962	0.17608	0.46	0.3207	-1.94926	-0.0023	1.47267
Residual Mass excess	2536	0.00179	0.41965	0.1761	0.46	0.32069	-1.94846	-0.0014	1.47305
Residual Atomic mass	2536	-1.04E-05	0.41962	0.17608	0.46	0.3207	-1.94926	-0.0023	1.47267

Where the abbreviations are as follows: SD = Standard Deviation, RMS = Root Mean Square and MAD = Mean Absolute Deviation.

Table 2. Comparison experimental and calculated data of the elements, which have a residual energy greater than module 1.5 MeV.

El Name	A	Z	N	Be	ResBe	AtMass	ResAtMass
Ca	53	20	33	8.33	3.69 E-02	49339.795	-1.96
Pb	208	82	126	7.87	9.04 E-03	193729.006	-1.89
Sn	132	50	82	8.35	1.23 E-02	122880.663	-1.63
Nh	278	113	165	7.18	5.74 E-03	259114.241	-1.52
No	259	102	157	7.40	5.80 E-03	241351.028	-1.51

Where the abbreviations are as follows: Be – Experimental Binding energy; ResBe – Difference of the calculated and experimental binding energy; AtMas – Experimental Atomic mass

ResAtMass- Difference of the calculated and experimental atomic mass

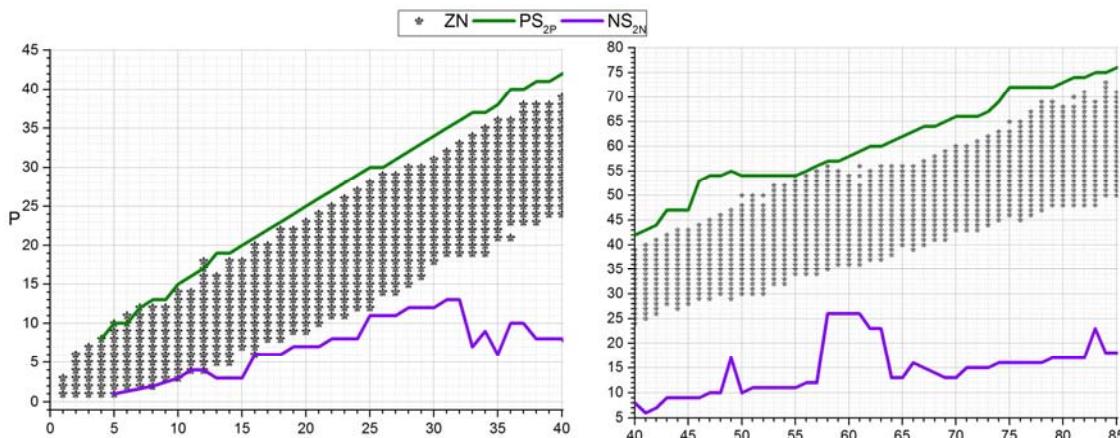
## 5. Proton and Neutron Drip Lines and Predictions

The definition of two proton and two neutron drip lines as a boundary of existing nuclear matter is as follow:

$$(Z+N) B_E(Z, N, a) > (Z+N-2) B_E(Z-2, N, a), \text{ for protons} \quad (5)$$

and

$$(Z+N) B_E(Z, N, a) > (Z+N-2) B_E(Z, N-2, a), \text{ for neutrons} \quad (6)$$



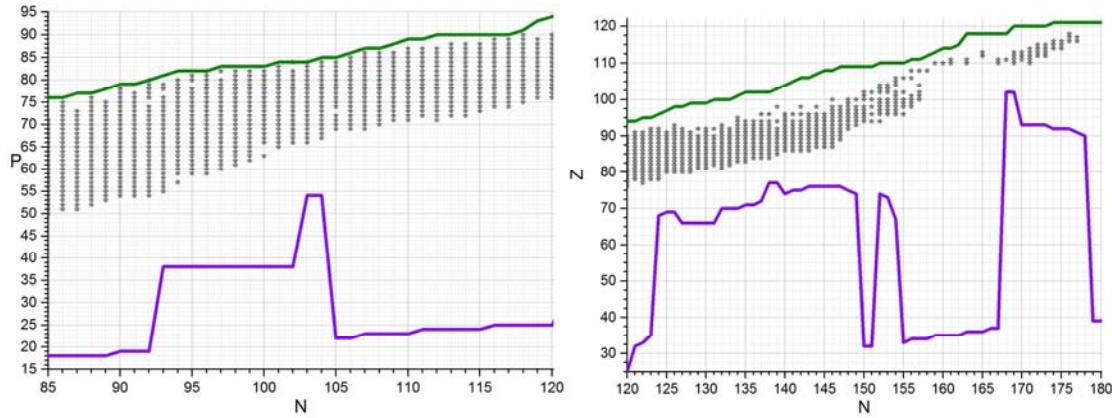
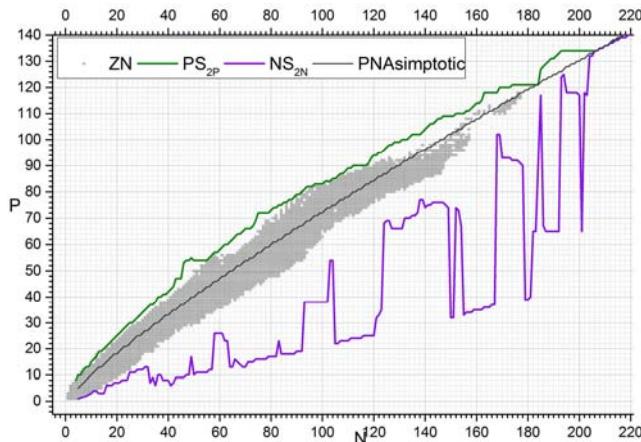
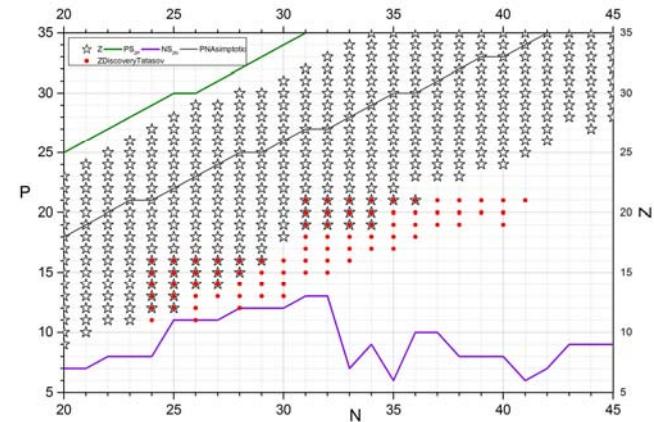


Figure 9. Two proton and neutron drip lines.

Figure 10. Asymptotic of two proton and neutron drip lines-  $Z = 134$ ,  $N = 206$ .Figure 11. The comparison of our model with the  $^{60}\text{Ca}_{20}$  RIKEN experiment [23] (PHYSICAL REVIEW LETTERS 121, 022501 (2018)).Table 3. Presents the Atomic Masses and Nuclear Binding Energies [MeV] for comparison with the data for discovered isotopes from the  $^{60}\text{Ca}_{20}$  RIKEN experiment [23].

No	El. Name	A	Z	N	AtMass	NBe
1	Sc	62	21	41	57755.72	480.90
2	Sc	61	21	40	56818.95	478.11
3	Sc	60	21	39	55884.56	472.93
4	Sc	59	21	38	54944.07	473.85
5	Sc	58	21	37	54010.17	468.19
6	Sc	57	21	36	53072.85	465.95
7	Sc	56	21	35	52138.51	460.72
8	Sc	55	21	34	51201.14	458.53
9	Sc	54	21	33	50267.14	452.96
10	Sc	53	21	32	49330.48	450.05
11	Sc	52	21	31	48397.49	443.48
12	Ca	60	20	40	55893.50	464.77
13	Ca	59	20	39	54958.81	459.89
14	Ca	58	20	38	54018.65	460.50
15	Ca	57	20	37	53085.43	454.15
16	Ca	56	20	36	52148.37	451.64
17	Ca	55	20	35	51213.94	446.50
18	Ca	54	20	34	50276.19	444.69
19	Ca	53	20	33	49341.74	439.57
20	Ca	52	20	32	48404.40	437.35
21	Ca	51	20	31	47470.86	431.32
22	K	59	19	40	54966.48	453.01
23	K	57	19	38	53095.12	445.24

No	El. Name	A	Z	N	AtMass	NBe
24	K	56	19	37	52163.60	437.19
25	K	55	19	36	51227.39	433.84
26	K	54	19	35	50293.77	427.90
27	K	53	19	34	49355.96	426.13
28	K	52	19	33	48421.41	421.12
29	K	51	19	32	47483.59	419.38
30	K	50	19	31	46549.55	413.85
31	Ar	54	18	36	50302.71	419.73
32	Ar	53	18	35	49369.29	413.59
33	Ar	52	18	34	48433.27	410.05
34	Ar	51	18	33	47498.22	405.53
35	Ar	50	18	32	46561.05	403.13
36	Ar	49	18	31	45625.96	398.66
37	Cl	52	17	35	48447.80	396.30
38	Cl	51	17	34	47512.47	392.06
39	Cl	50	17	33	46578.79	386.17
40	Cl	49	17	32	45641.77	383.63
41	Cl	48	17	31	44706.83	379.01
42	S	49	16	33	45655.56	370.62
43	S	48	16	32	44719.28	367.34
44	S	47	16	31	43785.70	361.35
45	S	46	16	30	42848.27	359.22
46	S	45	16	29	41913.87	354.05
47	S	44	16	28	40976.20	352.15
48	S	43	16	27	40042.19	346.60
49	S	42	16	26	39105.51	343.72
50	S	41	16	25	38172.97	336.69
51	S	40	16	24	37237.38	332.71
52	p	47	15	32	43796.90	350.93
53	p	46	15	31	42864.37	343.90
54	p	45	15	30	41928.20	340.51
55	p	44	15	29	40994.83	334.30
56	p	43	15	28	40058.12	331.45
57	p	42	15	27	39123.64	326.36
58	p	41	15	26	38185.58	324.86
59	p	40	15	25	37251.61	319.27
60	p	39	15	24	36314.73	316.58
61	Si	44	14	30	41007.40	322.53
62	Si	43	14	29	40073.00	317.35
63	Si	42	14	28	39136.70	314.09
64	Si	41	14	27	38203.07	308.16
65	Si	40	14	26	37265.56	306.10
66	Si	39	14	25	36331.28	300.81
67	Si	38	14	24	35393.10	299.43
68	Al	43	13	30	40098.72	292.41
69	Al	42	13	29	39160.50	291.08
70	Al	41	13	28	38220.57	291.44
71	Al	40	13	27	37284.99	287.46
72	Al	39	13	26	36347.01	285.86
73	Al	38	13	25	35413.08	280.23
74	Al	37	13	24	34475.37	278.38
75	Mg	40	12	28	37309.14	264.09
76	Mg	38	12	26	35429.05	265.04
77	Mg	37	12	25	34493.25	261.28
78	Mg	36	12	24	33554.40	260.56
79	Na	37	11	26	34521.74	233.57
80	Na	35	11	24	32641.30	234.88

In the following figure (Figure 12) we present the ZN coordinates, limited from the drip lines.

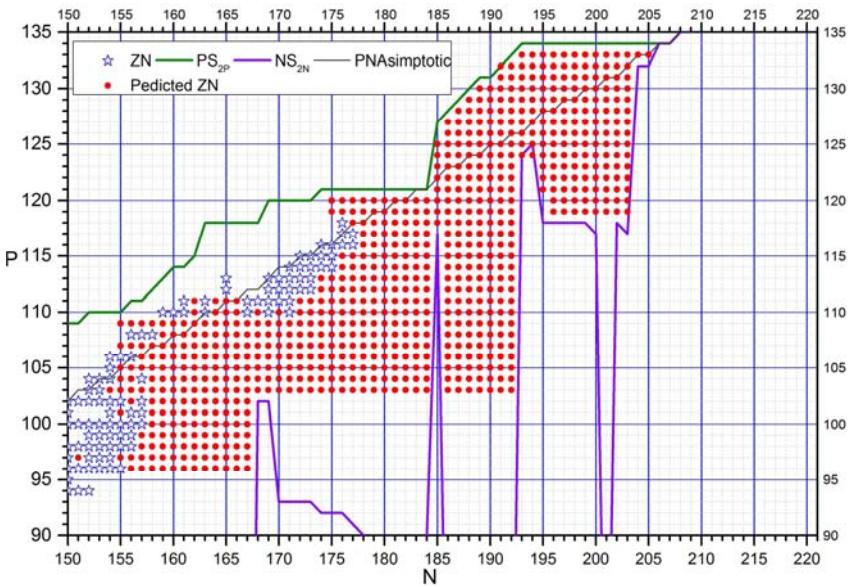


Figure 12. The ZN coordinates for which we will present the values of the Nuclear Binding Energy and Atomic Mass [MeV].

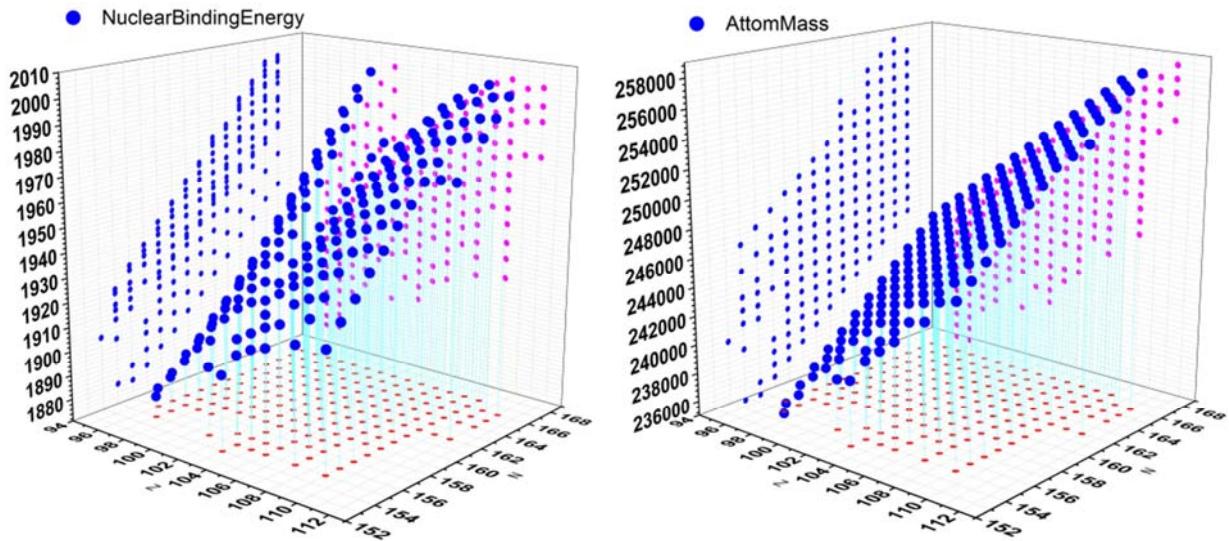


Figure 13. The values of Predicted Nuclear Binding Energies and Atomic Masses [MeV], Z in the interval 98-110 and N in the interval 154-167.

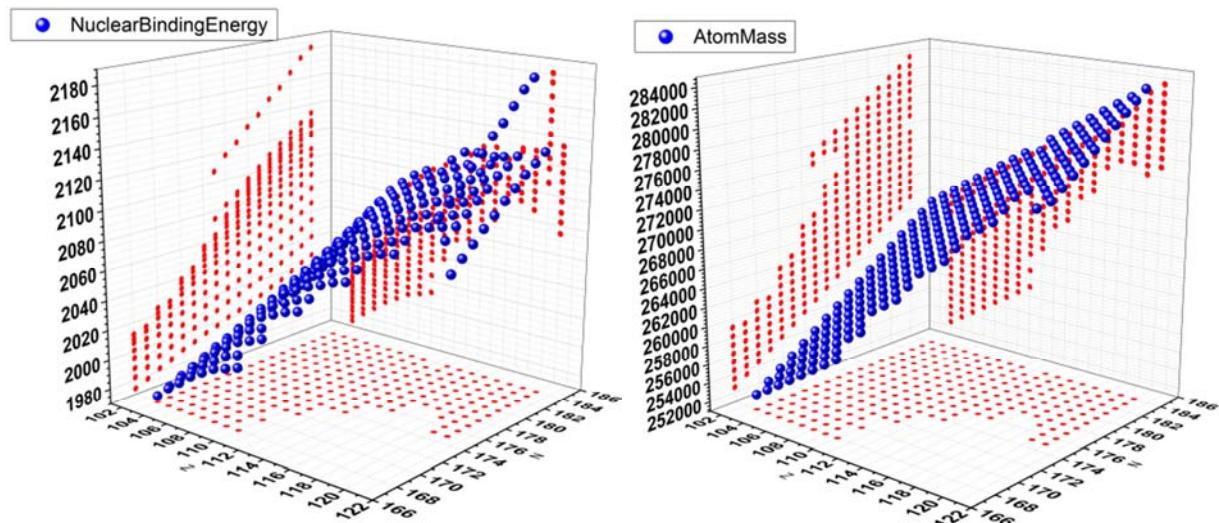
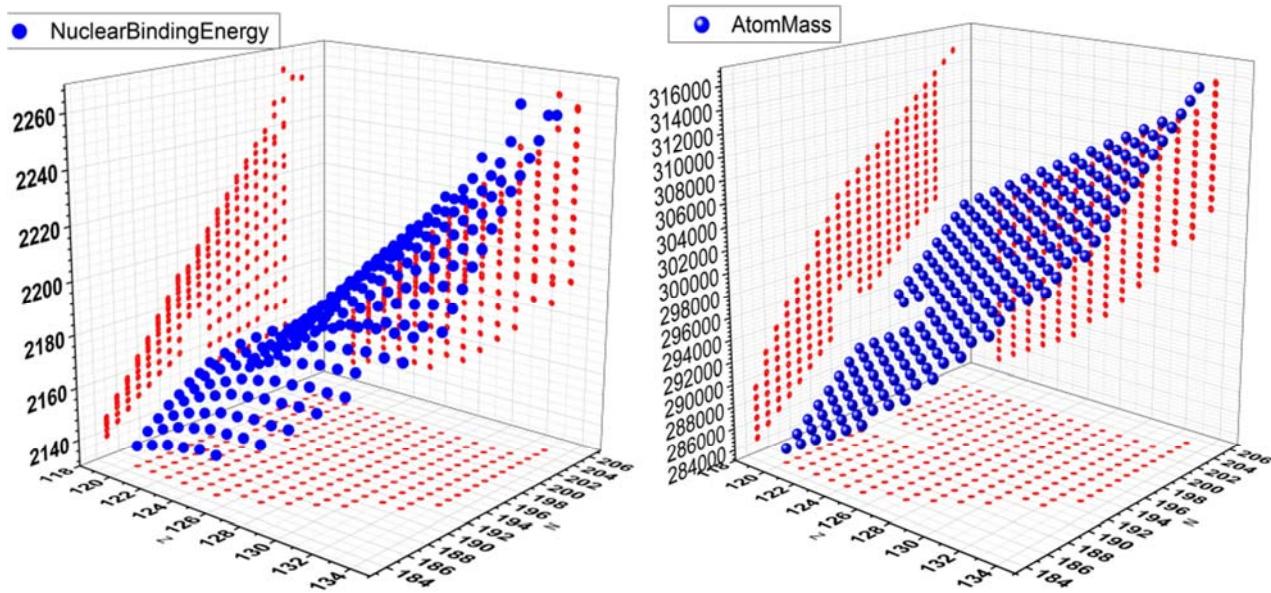
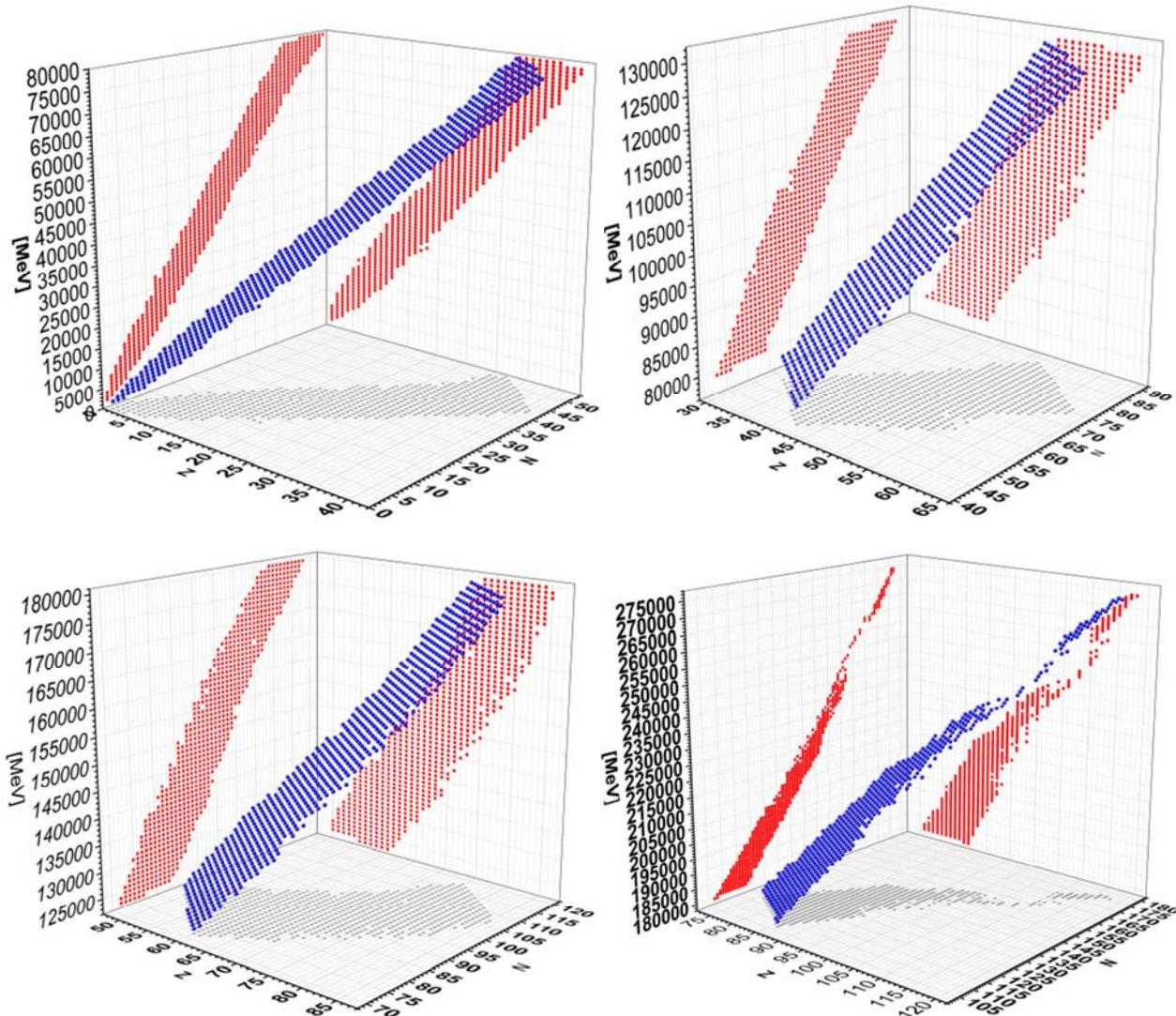


Figure 14. The values of Predicted Nuclear Binding Energies and Atomic Masses [MeV], Z in the interval 102-120 and N in the interval 168-184.



**Figure 15.** The values of Predicted Nuclear Binding Energies and Atomic Masses [MeV] Z in the interval 120-133 and N in the interval 185-205.



**Figure 16.** Illustration of the linear behavior of atomic masses in the Z, N plane, discovered in the Fig. 13- 15, for all the elements in the Mendeleev table.

**Table 4.** Presents the values of elements 119 and 120 atomic masses compared with the values calculated from our model and their residuals in MeV. The lines in bold present the data for predicted isotopes.

No	Element Name	A	Z	N	AtM	ResAtM
1	100Fm155	255	100	155	237614.773	-0.468
2	100Fm156	256	100	156	247862.541	
3	101Md156	257	101	156	239482.954	-0.323
4	101Md158	259	101	158	242285.848	
5	102No152	254	102	152	236684.203	-0.132
6	102No158	260	102	158	242285.848	
7	103Lr152	255	103	152	237620.921	0.615
8	102Lr154	257	103	154	239486.085	
9	104Rf157	261	104	157	243221.256	-0.344
10	104Rf158	262	104	158	244154.555	
11	105Db154	259	105	154	241358.940	0.683
12	105Db155	260	105	155	242291.383	
13	106Sg156	262	106	156	244159.797	0.679
14	106Sg157	263	106	157	245092.608	
15	107?156	263	107	156	245096.059	
16	107?157	264	107	157	246029.365	
17	108Hs158	266	108	158	247898.544	0.852
18	108Hs159	267	108	159	248830.841	
19	109?158	267	109	158	248836.155	
20	109?159	268	109	159	249769.155	
21	110Ds171	281	110	171	261903.072	1.253
22	110Gs172	282	110	172	262834.841	
23	111RG171	282	111	171	262838.858	0.666
24	111Rg172	283	111	172	263771.015	
25	112CN173	285	112	173	265640.791	0.659
26	112Cn174	286	112	174	266572.917	
27	113Nh173	286	113	173	266577.027	0.348
28	113Nh174	287	113	174	267509.317	
29	114FL175	289	114	175	269379.158	0.463
30	114Pl176	290	114	176	270311.375	
31	115Mc175	290	115	175	270315.827	0.523
32	115Mc176	291	115	176	271248.301	
33	116LV177	293	116	177	273118.238	-0.065
34	178Lv178	294	116	178	274051.565	
35	117Ts177	294	117	177	274055.298	0.001
36	117Tb178	295	117	178	274988.105	
37	118Og176	294	118	176	274058.520	-0.139
38	118Og177	295	118	177	274991.989	
39	118Og178	296	118	178	275924.647	
40	119?176	296	119	176	274961.336	
41	119?177	296	119	177	275893.923	
42	119?178	297	119	178	276824.007	
43	119?179	298	119	179	277756.839	
44	120?176	296	120	176	275934.877	
45	120?177	297	120	177	276867.829	
46	120?178	298	120	178	277799.928	
47	120?179	299	120	179	278733.174	

## 6. Discussion

Ormula terms describing the dependence of electron shells on nuclear mass.

The linear independence of the used arguments and value of their modules always were less than 1, provided simplification of calculations. The solution of the overdetermined system of

nonlinear equations has been obtained with the help of the Lubomir Aleksandrov's auto-regularization method of the Gauss-Newton type for ill-posed problems.

The impact of the electron shell was determined as a nonmonotonic parabolic function which is equal to zero at the magic number of electrons ( $E=Em$ ).

The result of the calculations allowed improved accuracy of nuclei mass estimation from 3.5 MeV of the original formula

and 2.2MeV [17] to 1.8 MeV using an improved numerical generalization of the Bethe-Weizsäcker Mass Formula.

The development of this work may include a calculation of the full and kinetic energy of decay as well as an estimation of nuclei lifetime.

## 7. Conclusion

The generalization of the Bethe –Weizsäcker formulae was approached by insertion of additional terms estimating the influence of proton, neutron and electron magic numbers on the nuclear mass and binding energy. This approach allows the describing of atomic masses starting from  $2\text{H}1$  to  $294\text{Og}118$  with RMS=0.46 MeV and with residuals of less than or equal to 1.9 MeV.

For only five elements the residuals are greater than 1.5 MeV.

The resulting agreement with the experimental data permits us to calculate realistic two proton and neutron drip lines with an asymptotic point at  $Z=132$ ,  $N=212$ .

There is an accordance of received model of nuclear and atomic masses with RICKEN experiment for creating of new  $20\text{Ca}20$  Isotopes [20].

The prediction for the masses of  $295\text{Og}118$  ( $118?177$ ),  $296\text{Og}118$  ( $118?178$ ),  $119?176$ ,  $120?176$  and some their isotopes are presented.

The like linear dependences discovered in  $Z$ ,  $N$  plane for all elements in Mendeleev table are illustrated.

Recent published [3, 8-14, 16-20] and presented herein results allow the suggestion that the resonant interaction of the electron shell and the nuclei causes the existence of the magic numbers.

The existence of proton and neutron magic numbers is a result of the unknown strong nuclear interactions in the effective field of electron shell. Probably, the electron shell field effects the nucleon run length [24], nuclei volume, rate of decay, and difference of the nuclei shape from a sphere. Thus, the nuclei interaction with the electron shell may switch on/off the strong interaction between neutrons and protons.

The other possible mechanism of such resonant interaction is inner oscillations of the protons and neutrons, which were proposed by M. Gryzinski (1959) [24] and N. Chetaev (1931-1962) [25]. The conjugated oscillations of the neutrons and protons may lead to Hopf (Poincaré-Andronov-Hopf) bifurcations in three-dimensional nuclear structures. Such bifurcations may correspond to periodic solutions and low frequency oscillations providing a breaking of nuclei symmetry and its interaction with electrons shells or decay.

The neutron and proton pair's spin (Gryzinski) interactions as proton oscillations are possible through nuclei - electrons resonance. The resonance and/or coherence may form an electronic bridge with resonant interactions between the nuclei of the atoms of most biological molecules, including H, C, N, P, etc. The totality of these molecules is the basis of the biosphere.

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